

# Homoclinic saddle to saddle-focus transitions in 3D and 4D ODEs

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Joint work with Manu Kalia & Hil Meijer



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## An example: Lorenz-Stenflo equations

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## General Methodology

Combine analytic and numerical  
bifurcation methods !

## Homoclinic orbits

- Consider a smooth family of smooth ODE systems

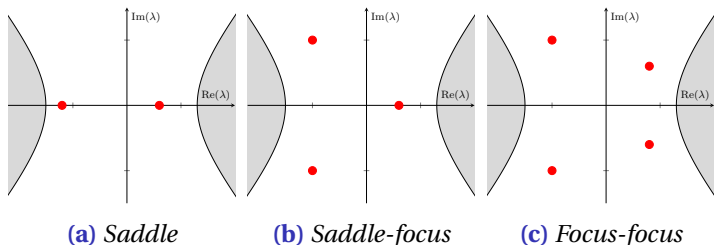
$$\dot{u} = F(u, p), \quad u \in \mathbb{R}^n, \quad p \in \mathbb{R}^m$$

- Suppose that at  $p_0 \in \mathbb{R}^m$  there is a *hyperbolic equilibrium*  $u_0 \in \mathbb{R}^n$ .
- Assume that  $u_0$  has a *primary homoclinic orbit*  $\Gamma_0$ , i.e. for a non-equilibrium solution  $t \mapsto u(t)$  holds

$$\lim_{t \rightarrow \pm\infty} u(t) = u_0$$

- Generically, the corresponding  $p_0$ -values occupy a codimension 1 homoclinic *bifurcation manifold* in  $\mathbb{R}^m$ .

- Leading eigenvalues of  $u_0$  (possibly, after reversing time):



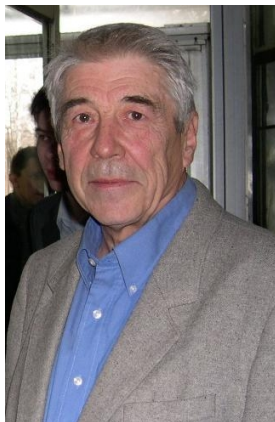
- Saddle quantity (index):

$$\sigma := \Re(\lambda^s) + \Re(\lambda^u) \quad \left( \nu := -\frac{\Re(\lambda^s)}{\Re(\lambda^u)} \right)$$

where  $\lambda^s$  and  $\lambda^u$  are leading stable and unstable eigenvalues of  $u_0$ .

## Limit cycles due to codim 1 homoclinic bifurcations

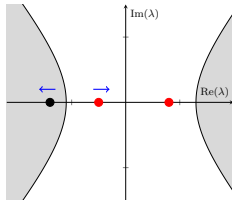
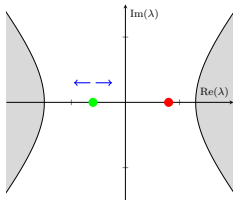
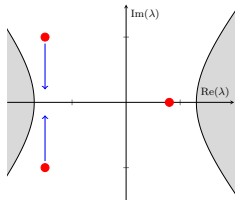
	$\sigma < 0$	$\sigma > 0$
<i>Saddle</i>	one cycle	one cycle
<i>Saddle-focus</i>	one cycle	$\infty$ cycles
<i>Focus-focus</i>	$\infty$ cycles	$\infty$ cycles



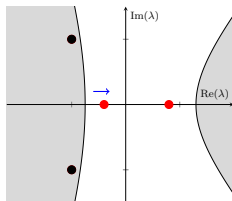
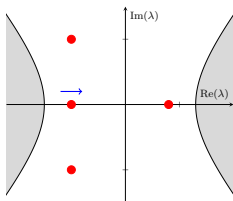
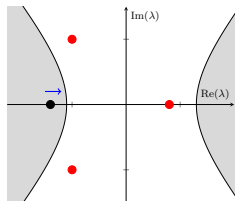
L.P. Shilnikov (1934-2011)

## Codim 2 saddle-focus to saddle transitions

2DL (double eigenvalue) Belyakov bifurcation:



3DL bifurcation:



$\mu_1 < 0$

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$\mu_1 = 0$

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$\mu_1 > 0$

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## 2DL Belyakov bifurcation



L.A. Belyakov

The bifurcation set in a system with a homoclinic saddle curve  
*Math. Notes* **28** (1980), 910-916



Yu.A. Kuznetsov, O. De Feo, and S. Rinaldi

Belyakov homoclinic bifurcations in a tritrophic food chain model  
*SIAM J. App. Math.* **62** (2001), 462-487

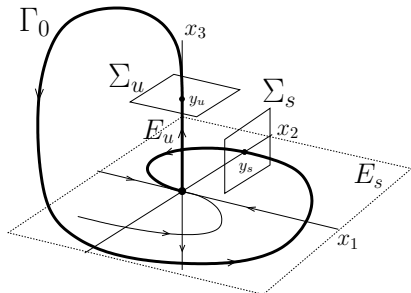
3D ODE system near the equilibrium  $x = 0$ :

$$\begin{cases} \dot{x}_1 &= \gamma(\mu)x_1 + x_2 + f_1(x, \mu) \\ \dot{x}_2 &= -\mu_1 x_1 + \gamma(\mu)x_2 + f_2(x, \mu) \\ \dot{x}_3 &= \beta(\mu)x_3 \end{cases}$$

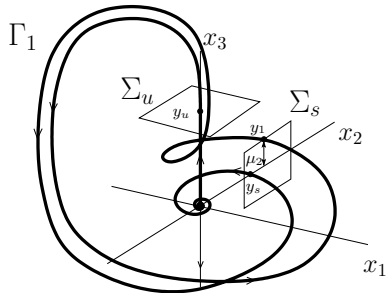
where  $\mu = (\mu_1, \mu_2) \in \mathbb{R}^2$  and  $\lambda_{1,2}^s = \gamma(0) < 0$ ,  $\lambda^u = \beta(0) > -\gamma(0)$ .



## Geometric setting



Homoclinic orbit at  $\mu = 0$



Double homoclinic orbit

## Model Poincaré maps

- 2D Model Poincaré Map on  $\Sigma_s$ :

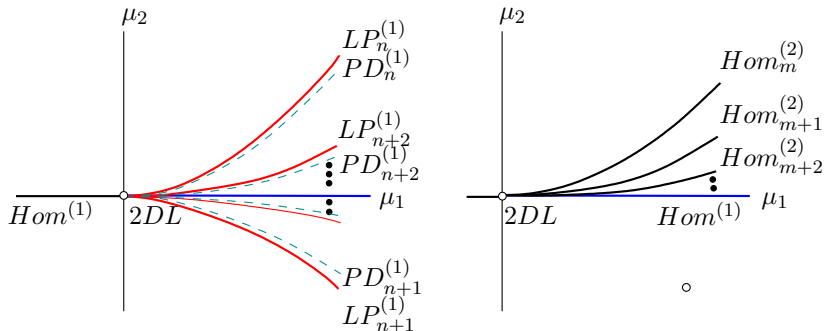
$$\begin{pmatrix} x_2 \\ x_3 \end{pmatrix} \mapsto \begin{pmatrix} 1 + \frac{A}{\sqrt{\mu_1}} x_2 x_3^v \sin\left(-\frac{\sqrt{\mu_1}}{\beta} \ln x_3 + \Theta\right) \\ \mu_2 + \frac{C}{\sqrt{\mu_1}} x_2 x_3^v \sin\left(-\frac{\sqrt{\mu_1}}{\beta} \ln x_3\right) \end{pmatrix}$$

where

$$v = -\frac{\gamma}{\beta} < 1 \quad (\text{chaotic case})$$

- 1D Model Poincaré Map:

$$x_3 \mapsto \mu_2 + \frac{C}{\sqrt{\mu_1}} x_3^v \sin\left(-\frac{\sqrt{\mu_1}}{\beta} \ln x_3\right)$$

**(Partial) bifurcation diagram ( $\nu < 1$ )**

$$LP_n^{(1)} \text{ and } PD_n^{(1)} : \mu_2^{(n)} = \frac{(-1)^n C}{e\gamma} \exp\left(\frac{\gamma\pi n}{\sqrt{\mu_1}}\right) (1 + o(\mu_1))$$

$$Hom_m^{(2)} : \mu_2^{(m)} = \exp\left(-\frac{\beta\pi m}{\sqrt{\mu_1}}\right) (1 + o(\mu_1))$$

## 3DL bifurcation

 M. Kalia, Yu.A. Kuznetsov, and H.G.E. Meijer

Homoclinic saddle to saddle-focus transitions in 4D systems

[arXiv:1712.03212](https://arxiv.org/abs/1712.03212)[submitted to Nonlinearity]

4D ODE  $C^1$  orbitally equivalent system near the equilibrium  $x = 0$ :

$$\begin{cases} \dot{x}_1 &= \gamma(\mu)x_1 - x_2, \\ \dot{x}_2 &= x_1 + \gamma(\mu)x_2, \\ \dot{x}_3 &= (\gamma(\mu) - \mu_1)x_3, \\ \dot{x}_4 &= \beta(\mu)x_4, \end{cases}$$

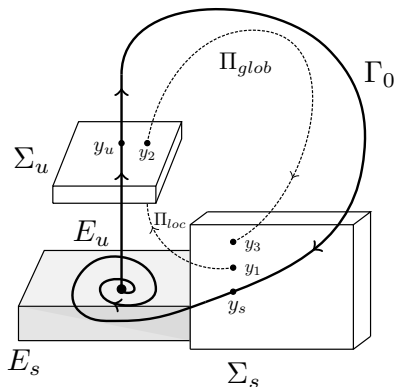
where  $\mu = (\mu_1, \mu_2) \in \mathbb{R}^2$  and

$$\Re(\lambda_{1,2,3}^s) = \gamma(0) < 0, \quad \lambda^u = \beta(0) > -\gamma(0)$$

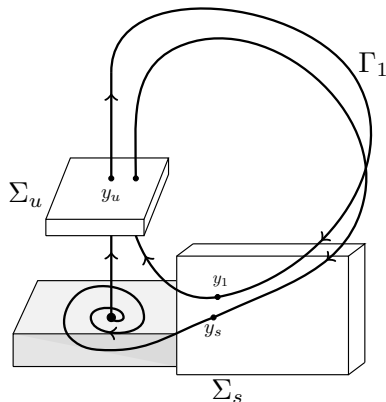
so that

$$\nu = -\frac{\gamma}{\beta} < 1 \quad (\text{chaotic case})$$

## Geometric setting



Homoclinic orbit



Double homoclinic orbit

## Model Poincaré maps

- 3D Model Poincaré Map on  $\Sigma_s$ :

$$G: \begin{pmatrix} x_1 \\ x_2 \\ x_4 \end{pmatrix} \mapsto \begin{pmatrix} 1 + \alpha_1 x_1 x_4^v \cos\left(-\frac{1}{\beta} \ln x_4 + \phi_1\right) + \alpha_2 x_2 x_4^{v+\mu_1/\beta} \\ 1 + \alpha_3 x_1 x_4^v \sin\left(-\frac{1}{\beta} \ln x_4 + \phi_2\right) + \alpha_4 x_2 x_4^{v+\mu_1/\beta} \\ \mu_2 + C_1 x_1 x_4^v \sin\left(-\frac{1}{\beta} \ln x_4\right) + C_2 x_2 x_4^{v+\mu_1/\beta} \end{pmatrix}$$

- 1D Model Poincaré Map:

$$F: x_4 \mapsto \mu_2 + C_1 x_4^v \sin\left(-\frac{1}{\beta} \ln x_4\right) + C_2 x_4^{v+\mu_1/\beta}$$

## Lemma (Bifurcations of 1D model map)

Generically, the model map  $F$  with  $\nu < 1$  has an infinite number of fold curves  $LP_n^{(1)}$  accumulating to the half axis  $\mu_2 = 0$  with  $\mu_1 \geq 0$ . Each curve resembles a 'horn' with a cusp point

$$CP_n^{(1)} : \begin{pmatrix} \mu_1^{(n)} \\ \mu_2^{(n)} \end{pmatrix} = \begin{pmatrix} \frac{1}{4\pi n} [\ln a + O(\frac{1}{n})] \\ -e^{-\beta\nu(2\pi n + \theta + \phi)} \frac{\text{sign}(C_2)C_1}{\beta\nu\sqrt{1+\beta^2\nu^2}} a^{-(\theta+\phi)/4\pi n} + O(\frac{1}{\sqrt{n}}) \end{pmatrix}$$

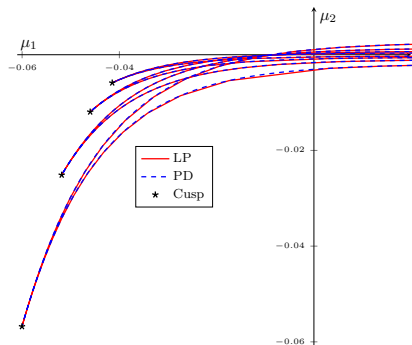
where

$$a = \frac{\beta^2\nu^2}{1 + \beta^2\nu^2} \frac{C_2^2}{C_1^2}, \quad \sin\phi = (1 + \beta^2\nu^2)^{-1/2}, \quad \theta = \begin{cases} \pi/2, & \text{if } C_2 < 0, \\ 3\pi/2, & \text{if } C_2 > 0. \end{cases}$$

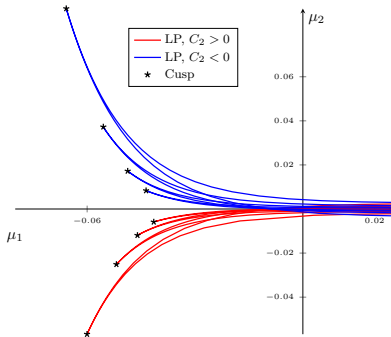
Moreover, there exists an infinite number of period-doubling curves  $PD_n^{(1)}$  having – away from the cusp points  $CP_n^{(1)}$  – the same asymptotic properties as the fold bifurcation curves  $LP_n^{(1)}$ .

# Primary LP and PD curves of 1D map $F$

$$C_1 = 1.2, C_2 = 0.7$$



$$C_1 = 1.2, C_2 = \pm 0.7$$



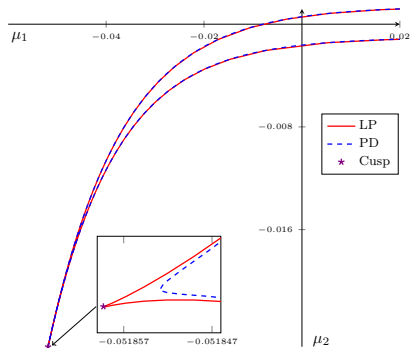
$$\beta = 0.2, \quad \nu = 0.5$$



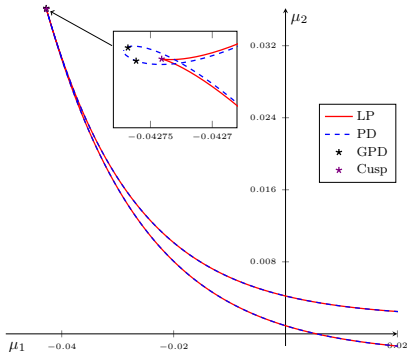
## Spring and saddle areas

Depending on  $(C_1, C_2)$ , curves  $PD_n^{(1)}$  could either be smooth or develop small loops around the corresponding cusp points

$C_1 = 1.2, C_2 = 0.7$  (Saddle area)



$C_1 = 0.8, C_2 = -1.1$  (Spring area)



## Theorem (Bifurcations of 3D model map)

*Generically, the model map  $G$  with  $\nu < 1$  has an infinite number of fold and period-doubling curves accumulating to the half axis  $\mu_2 = 0$  with  $\mu_1 \geq 0$  and having the same asymptotic properties as the curves  $LP_n^{(1)}$  and  $PD_n^{(1)}$  above.*

*Moreover, there exists an infinite sequence of ‘parabolas’  $\text{Hom}_m^{(2)}$  that accumulate onto the half axis  $\mu_2 = 0$  with  $\mu_1 \geq 0$  and correspond to the double homoclinic orbits. The turning points of the parabolas are*

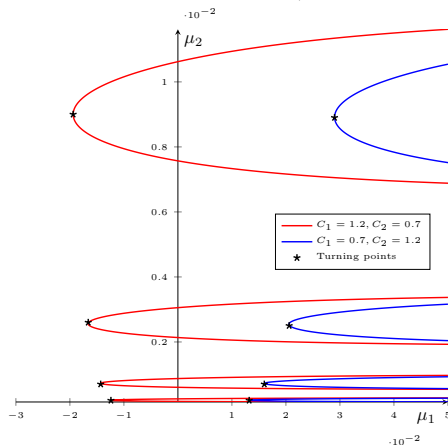
$$\begin{pmatrix} \mu_1^{(m)} \\ \mu_2^{(m)} \end{pmatrix} = \begin{pmatrix} \frac{1}{4\pi m} \left( \ln \left( \frac{C_2^2}{C_1^2} \right) + O\left(\frac{1}{m}\right) \right) \\ e^{-\beta(2\pi m + \theta)} \left( 1 + O\left(\frac{1}{m}\right) \right) \end{pmatrix}$$

where

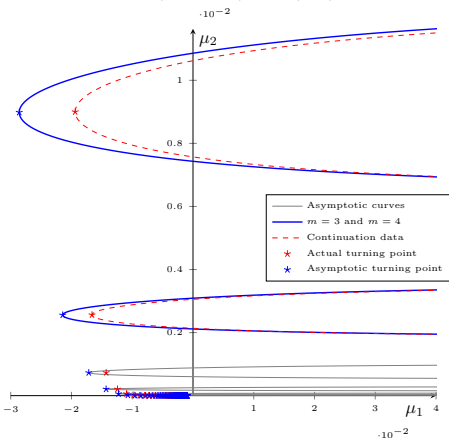
$$\theta = \begin{cases} \pi/2, & \text{if } C_2 < 0, \\ 3\pi/2, & \text{if } C_2 > 0. \end{cases}$$

# Secondary homoclinic curves defined by 3D map $G$

Continuations:  $m = 3, 4 \dots 6$



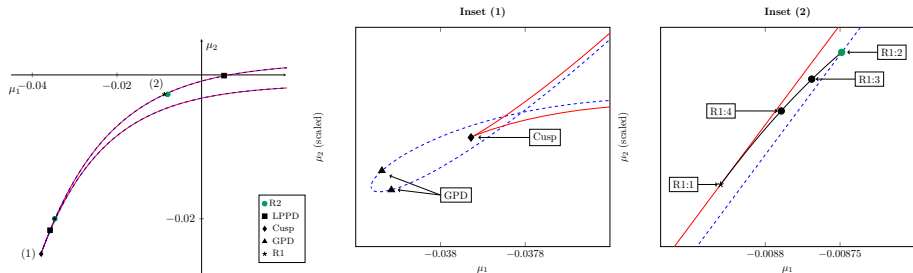
$C_1 = 1.2, C_2 = 0.7, m = 3, 4 \dots, 90$



$$\beta = 0.2, \quad \nu = 0.5$$

# Codim 2 bifurcations in 3D map $G$

$$\nu = 0.5, \beta = 0.2$$



$$C_1 = 0.8, C_2 = 1.2, \alpha_1 = 0.8, \alpha_2 = 1.3, \alpha_3 = 0.6, \alpha_4 = 1.1$$

$$\phi_1 = \phi_2 = \pi/6$$



## An example: Modified Lorenz-Stenflo equations

L. Stenflo

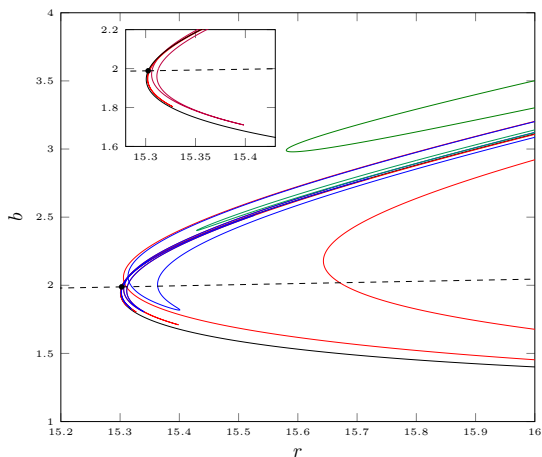
Generalized Lorenz equations for acoustic-gravity waves in the atmosphere

*Physica Scripta* **53**(1):83, 1996

$$\begin{cases} \dot{x} &= \sigma(y-x) + su \\ \dot{y} &= rx - xz - y + \epsilon_1 z \\ \dot{z} &= xy - bz \\ \dot{u} &= -x - \sigma u + \epsilon_2 y \end{cases}$$

$$\sigma = 0.1, s = 33, \epsilon_1 = 0.1, \epsilon_2 = 0.3$$

## Chaotic 3DL bifurcation in modified LS model



$$3DL: (r, b) \approx (15.302531, 1.9884) \quad (\nu \approx 0.71605 < 1)$$

## Open questions

- Exact condition for the Saddle/Spring area transition in 2D model map.
- Analytic proof of strong resonances in 3D map.
- Can  $C^k$ -linearization with  $k > 1$  and related non-resonance conditions be avoided ?
- Can the use of Homoclinic Center Manifold Theorem and related gap conditions be avoided as well ?
- More realistic examples.
- What happens in the volume-preserving case, where 3DL has codim 1 and is always chaotic ?