

22c. Case $0 < m_0(j_2) < p$, and $\varepsilon = -1$

```
In[ = ]:= Clear[r0]
F = tht[0] f[r0, t] Phi[h, p, r0, p] + tht[1] f[r0 - 2, t] Phi[h, p, r0 - 2, p]
```

```
Out[ = ]:= f[r0, t] Phi[h, p, r0, p] tht[0] + f[-2 + r0, t] Phi[h, p, -2 + r0, p] tht[1]
```

We use that one shift operator gives zero:

```
In[ = ]:= sh[-3, -1, F, subnab] /. eps → -1 // Simplify
Coefficient[%, Phi[-3 + h, -1 + p, -1 + r0, -1 + p]]
fr0m2 = f[r0 - 2, t] /. Solve[% == 0, f[r0 - 2, t]][1] // Simplify

Out[ = ]:= - $\frac{1}{4 \times (1 + p)}$  p (4 i  $\sqrt{2 \pi}$  t  $\sqrt{\text{Abs}[ell]}$  f[-2 + r0, t] Phi[-3 + h, -1 + p, -1 + r0, -1 + p] tht[0] +
(4 + h + 2 p - r0 + 4 ell  $\pi$  t2) f[r0, t] Phi[-3 + h, -1 + p, -1 + r0, -1 + p] tht[0] +
6 f[-2 + r0, t] Phi[-3 + h, -1 + p, -3 + r0, -1 + p] tht[1] +
h f[-2 + r0, t] Phi[-3 + h, -1 + p, -3 + r0, -1 + p] tht[1] +
2 p f[-2 + r0, t] Phi[-3 + h, -1 + p, -3 + r0, -1 + p] tht[1] -
r0 f[-2 + r0, t] Phi[-3 + h, -1 + p, -3 + r0, -1 + p] tht[1] +
4 ell  $\pi$  t2 f[-2 + r0, t] Phi[-3 + h, -1 + p, -3 + r0, -1 + p] tht[1] -
2 t Phi[-3 + h, -1 + p, -3 + r0, -1 + p] tht[1] f^{(0,1)}[-2 + r0, t] -
2 t Phi[-3 + h, -1 + p, -1 + r0, -1 + p] tht[0] f^{(0,1)}[r0, t])
```



```
Out[ = ]:= - $\frac{1}{4 \times (1 + p)}$  p (4 i  $\sqrt{2 \pi}$  t  $\sqrt{\text{Abs}[ell]}$  f[-2 + r0, t] tht[0] +
(4 + h + 2 p - r0 + 4 ell  $\pi$  t2) f[r0, t] tht[0] - 2 t tht[0] f^{(0,1)}[r0, t])
```



```
Out[ = ]:=  $\frac{i ((4 + h + 2 p - r0 + 4 ell \pi t^2) f[r0, t] - 2 t f^{(0,1)}[r0, t])}{4 \sqrt{2 \pi} t \sqrt{\text{Abs}[ell]}}$ 
```

Next we use the other shift operator

```
In[ 0]:= Fp = sh[3, -1, F, subnab] /. eps → -1 // Simplify
fp[r0 - 1, t] = Coefficient [%, Phi[3 + h, -1 + p, 1 + r0, -1 + p]] / tht[0] // Simplify
Out[ 0]= - $\frac{1}{4 \times (1 + p)}$ 
p (f[r0, t] ((4 - h + 2 p + r0 - 4 ell π t2) Phi[3 + h, -1 + p, 1 + r0, -1 + p] × tht[0] - 4 i  $\sqrt{2 \pi}$  t
 $\sqrt{\text{Abs}[ell]}$  Phi[3 + h, -1 + p, -1 + r0, -1 + p] × tht[1]) +
f[-2 + r0, t] ((2 - h + 2 p + r0 - 4 ell π t2) Phi[3 + h, -1 + p, -1 + r0, -1 + p] × tht[1] -
8 i  $\sqrt{\pi}$  t  $\sqrt{\text{Abs}[ell]}$  Phi[3 + h, -1 + p, -3 + r0, -1 + p] × tht[2]) -
2 t (Phi[3 + h, -1 + p, -1 + r0, -1 + p] × tht[1] f(0,1)[-2 + r0, t] +
Phi[3 + h, -1 + p, 1 + r0, -1 + p] × tht[0] f(0,1)[r0, t]))
Out[ 0]=  $\frac{p ((-4 + h - 2 p - r0 + 4 ell \pi t^2) f[r0, t] + 2 t f^{(0,1)}[r0, t])}{4 \times (1 + p)}$ 
```

We know the lowest order term of b_u^- .

Hence we have with a non-zero factor **cnz**:

```
In[ 0]:= Clear[cnz]
eq = ((-4 + h - 2 p - r0 + 4 ell π t2) f[r0, t] + 2 t f(0,1)[r0, t]) ==
cnz t^(0 + 1) WhittakerV[-p + 1 + (j2 - 1)/2, nu2, 2 Pi Abs[ell] t^2] /.
ell → -Abs[ell] /. hpsub // Simplify
Out[ 0]= (-4 + h - 2 p - r0 - 4 π t2 Abs[ell]) f[r0, t] + 2 t f(0,1)[r0, t] ==
cnz t WhittakerV [ $\frac{1}{4} \times (2 + h - p)$ ,  $\frac{1}{2} \text{Abs}[h - p]$ ,  $2 \pi t^2 \text{Abs}[ell]$ ]
```

We solve for the derivative of f_{r_0}

```
In[ 0]:= fr0d = f(0,1)[r0, t] /. Solve[eq, f(0,1)[r0, t]][1] // Simplify
Out[ 0]=  $\frac{1}{2 t} \left( (4 - h + 2 p + r0 + 4 \pi t^2 \text{Abs}[ell]) f[r0, t] +$ 
cnz t WhittakerV [ $\frac{1}{4} \times (2 + h - p)$ ,  $\frac{1}{2} \text{Abs}[h - p]$ ,  $2 \pi t^2 \text{Abs}[ell]$ ]  $\right)$ 
```

This can be inserted in the eigenfunction equations, together with the expressions for f_{r_0+2}

```
In[ = ei =
  efEqn[h, p, r0, f, ell, 0, -1] /. j → j2 /. nu → nu2 /. f[r0 - 2, t] → fr0m2 /. f^(0,1)[-2 + r0, t] →
  D[fr0m2, t] /. f^(0,2)[r0, t] → D[fr0d, t] /. f^(0,1)[r0, t] → fr0d // Whrel /.
  hpsub /. ell → -Abs[ell] /. Abs[h - p]^2 → (h - p)^2 // Simplify

Out[ = {1
  16 cnz t
  (-4 × (4 + 2 h - p + r0 - 8 π t^2 Abs[ell]) WhittakerV[(1
  4 × (2 + h - p), 1
  2 Abs[h - p], 2 π t^2 Abs[ell])] +
  (16 - 3 h^2 - 8 p - 3 p^2 + h (8 + 6 p)) WhittakerV[(1
  4 × (6 + h - p), 1
  2 Abs[h - p], 2 π t^2 Abs[ell])]),
  -3
  16 cnz (p + r0) t (4 × (4 + 2 h - p + r0 - 8 π t^2 Abs[ell])
  WhittakerV[(1
  4 × (2 + h - p), 1
  2 Abs[h - p], 2 π t^2 Abs[ell])] +
  (-16 + 3 h^2 + 8 p + 3 p^2 - 2 h (4 + 3 p)) WhittakerV[(1
  4 × (6 + h - p), 1
  2 Abs[h - p], 2 π t^2 Abs[ell])] )]
```

These expressions should be zero. We insert the asymptotic behavior of the V-Whittaker function.

```
In[ = Clear[tau]
  eia = ei /. t → Sqrt[tau]/Sqrt[2 Pi Abs[ell]] /. Abs[ell] → ell /.
  WhittakerV[kap_, s_, tau_] → -E^(-Pi I kap) tau^(-kap) E^(tau/2) // Simplify

Out[ = {1
  16 √ell √2 π i cnz e^(-1/4 i (h π - p π + 2 i tau)) tau^(1/4 (-4 - h + p))
  (3 h^2 + 3 p^2 + 4 p (2 + tau) - 2 h (4 + 3 p + 4 tau) - 4 × (4 + (4 + r0) tau - 4 tau^2)),
  1
  16 √ell √2 π 3 i cnz e^(-1/4 i (h π - p π + 2 i tau)) (p + r0) tau^(1/4 (-4 - h + p))
  (3 h^2 + 3 p^2 + 4 p (2 + tau) - 2 h (4 + 3 p + 4 tau) - 4 × (4 + (4 + r0) tau - 4 tau^2)) }]
```

```
In[ = eia E^(-tau/2) tau^(-1 + (h - p)/4) // Simplify
  Limit[% , tau → Infinity]

Out[ = { (i cnz e^(-1/4 i (h - p) π) (3 h^2 + 3 p^2 + 4 p (2 + tau) - 2 h (4 + 3 p + 4 tau) - 4 × (4 + (4 + r0) tau - 4 tau^2))) /
  (16 √ell √2 π tau^2), (3 i cnz e^(-1/4 i (h - p) π) (p + r0) (3 h^2 + 3 p^2 + 4 p (2 + tau) -
  2 h (4 + 3 p + 4 tau) - 4 × (4 + (4 + r0) tau - 4 tau^2))) / (16 √ell √2 π tau^2) }

Out[ = { i cnz e^(-1/4 i (h - p) π), e^(-1/4 i (h - p) π) (48 i cnz p + 48 i cnz r0)
  √ell √2 π, 16 √ell √2 π }
```

So **cnz** has to vanish.