

Turing Instability in a 1D Neural Wave Equation

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Synaptic Processing

- Post-synaptic

current: $I_s = g(V_s - V)$.

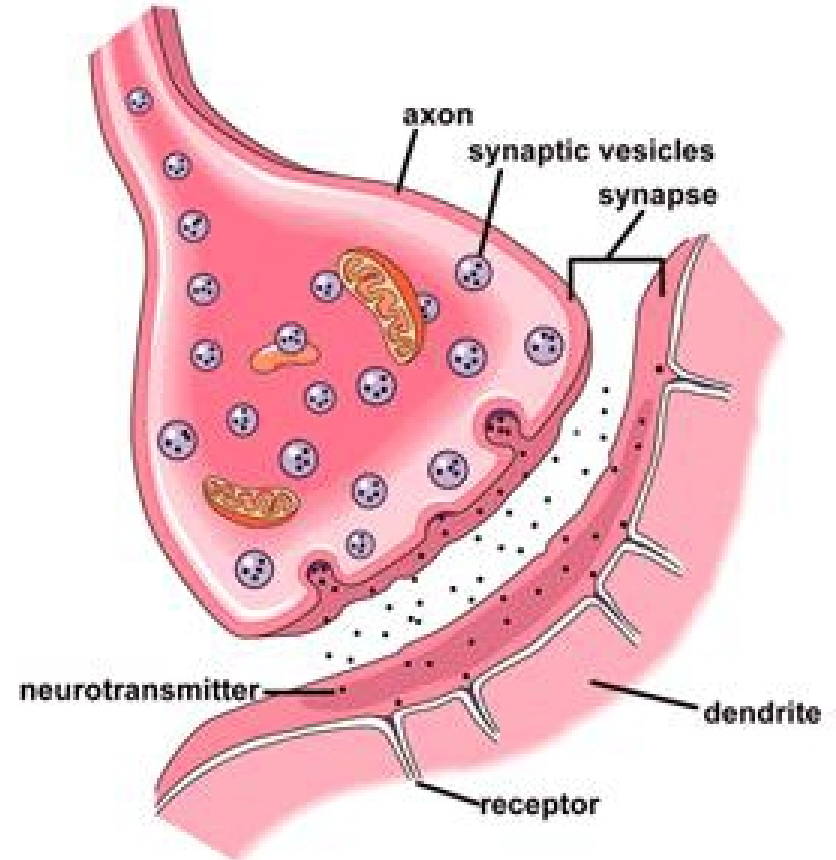
$V_s > 0$: excitatory,

$V_s < 0$: inhibitory.

$$g(t) = \bar{g}\eta(t - T), \quad t \geq T,$$

- Conductance

change: $g(t) = \bar{g} \sum_m \eta(t - T_m)$.



Synaptic Processing

- Two common choices for $\eta(t)$:

$$\eta(t) = \left(\frac{1}{\alpha} - \frac{1}{\beta} \right)^{-1} [e^{-\alpha t} - e^{-\beta t}] H(t),$$

$$\eta(t) = \alpha e^{-\alpha t} H(t).$$

- Conductance change from train of

APs:
$$g(t) = \bar{g} \sum_m \eta(t - T_m).$$

Dendritic Processing

Basic uniform cable equation:

$$\frac{\partial V(x, t)}{\partial t} = -\frac{V(x, t)}{\tau} + D \frac{\partial^2 V(x, t)}{\partial x^2} + I(x, t), \quad x \in (-\infty, \infty)$$

Green's function:

$$G_{\infty}(x, t) = \frac{1}{\sqrt{4\pi Dt}} e^{-t/\tau} e^{-x^2/(4Dt)}$$

General solution:

$$V(x, t) = \int_{-\infty}^t dt' \int_{-\infty}^{\infty} dx' G_{\infty}(x - x', t - t') I(x', t') + \int_{-\infty}^{\infty} dx' G_{\infty}(x - x', t) V(x', 0).$$

Firing Rates

- Spike train $Qg = \bar{g} \sum_m \delta(t - T_m)$.
- Short-time average: $Qg = f$, the instantaneous firing rate.
- For a single population with self-feedback we get equations like $Qg = w_0 f(g)$.

1D Tissue Level Model

$$Qg = \int_{-\infty}^{\infty} w(x, y) f(g(y, t - D(x, y)/v)) dy.$$

Voltage $V(\xi, x, t)$ at position $\xi \geq 0$ along cable:

$$\frac{\partial V}{\partial t} = -\frac{V}{\tau} + D \frac{\partial^2 V}{\partial \xi^2} + I(\xi, x, t).$$

I = synaptic input, proportional to a conductance change

$$g(\xi, x, t) = \int_{-\infty}^t ds \eta(t - s) \int_{-\infty}^{\infty} dy W(\xi, x, y) f(h(y, s - D(x, y)/v)).$$

1D Tissue Level Model

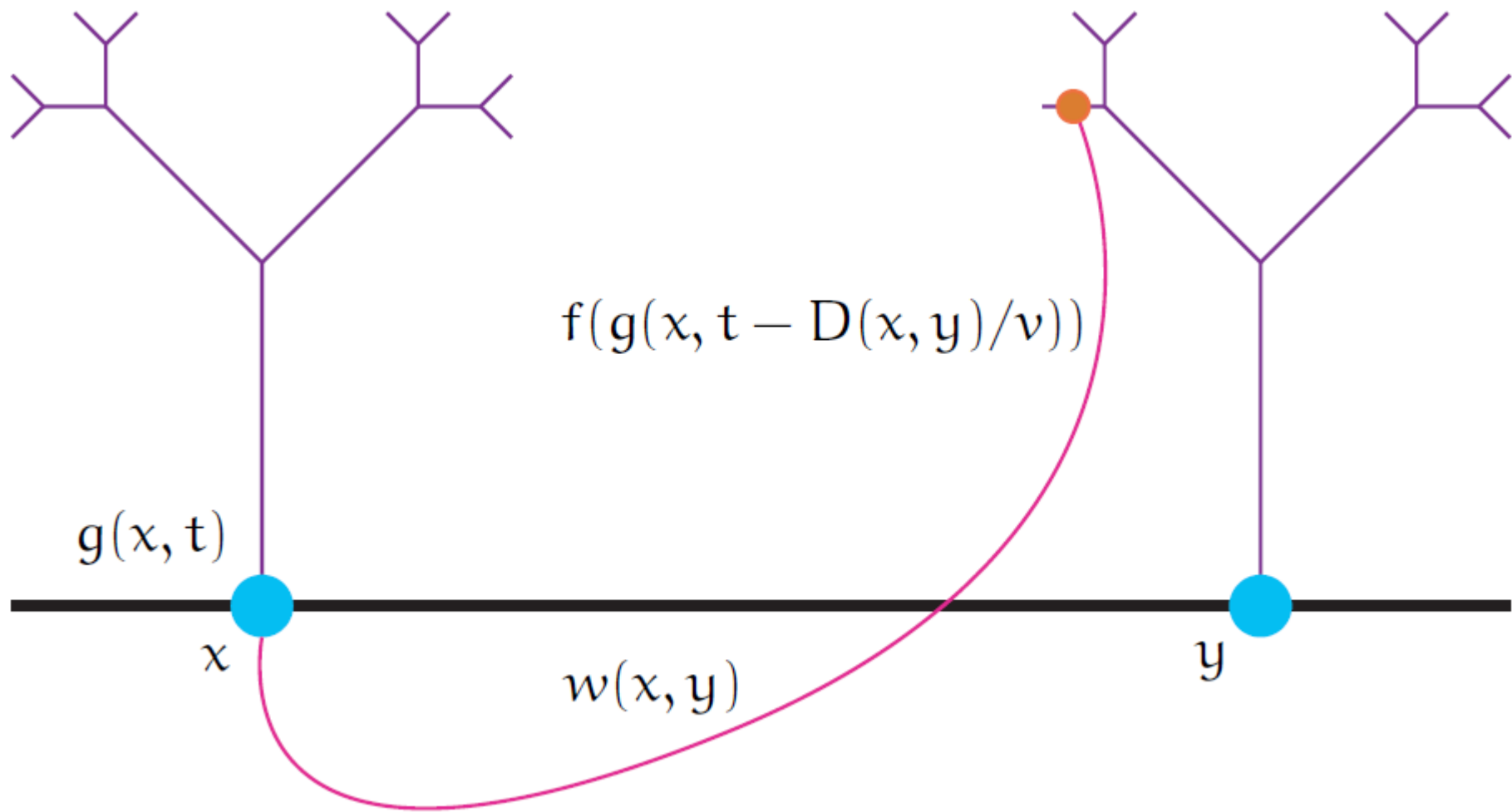
Axo-dendritic weights can be decomposed in the product form

$$W(\xi, x, y) = P(\xi)w(|x - y|).$$

Then the equation for h takes the form

$$h(x, t) = \kappa \int_{-\infty}^t ds F(t - s) \int_{-\infty}^s ds' \eta(s - s') \int_{-\infty}^{\infty} dy w(|x - y|) f(h(y, s' - D(x, y)/v))$$

$$F(t) = \int_0^{\infty} d\xi P(\xi) G(\xi, t).$$



One-dimensional model without dendrites or axonal

$$h(x, t) = \kappa \int_0^\infty \eta(s) \int_{-\infty}^\infty w(|y|) f(h(x - y, t - s)) dy ds.$$

Spatially uniform resting state $h(x, t) = h_0$, defined by

$$h_0 = \kappa f(h_0) \int_{-\infty}^\infty w(|y|) dy.$$

We linearize by letting $h(x, t) \rightarrow h_0 + h(x, t)$,

so that $f(h) \rightarrow f(h_0) + f'(h_0)h$.

Linearization around resting state

$$h(x, t) = \kappa \int_0^{\infty} \eta(s) \int_{-\infty}^{\infty} w(|y|) h(x - y, t - s) dy ds,$$

$$\beta = f'(h_0).$$

Solutions are of the form $e^{\lambda t} e^{ipx}$, with

$$1 = \kappa \beta \tilde{\eta}(\lambda) \hat{w}(p), \quad \hat{w}(p) = \int_{-\infty}^{\infty} w(|y|) e^{-ipy} dy,$$

$$\tilde{\eta}(\lambda) = \int_0^{\infty} \eta(s) e^{-\lambda s} ds.$$

Instability Analysis

The uniform steady state is linearly stable if

$$\operatorname{Re}(\lambda(p)) < 0 \quad \forall p \in \mathbb{R} \setminus 0.$$

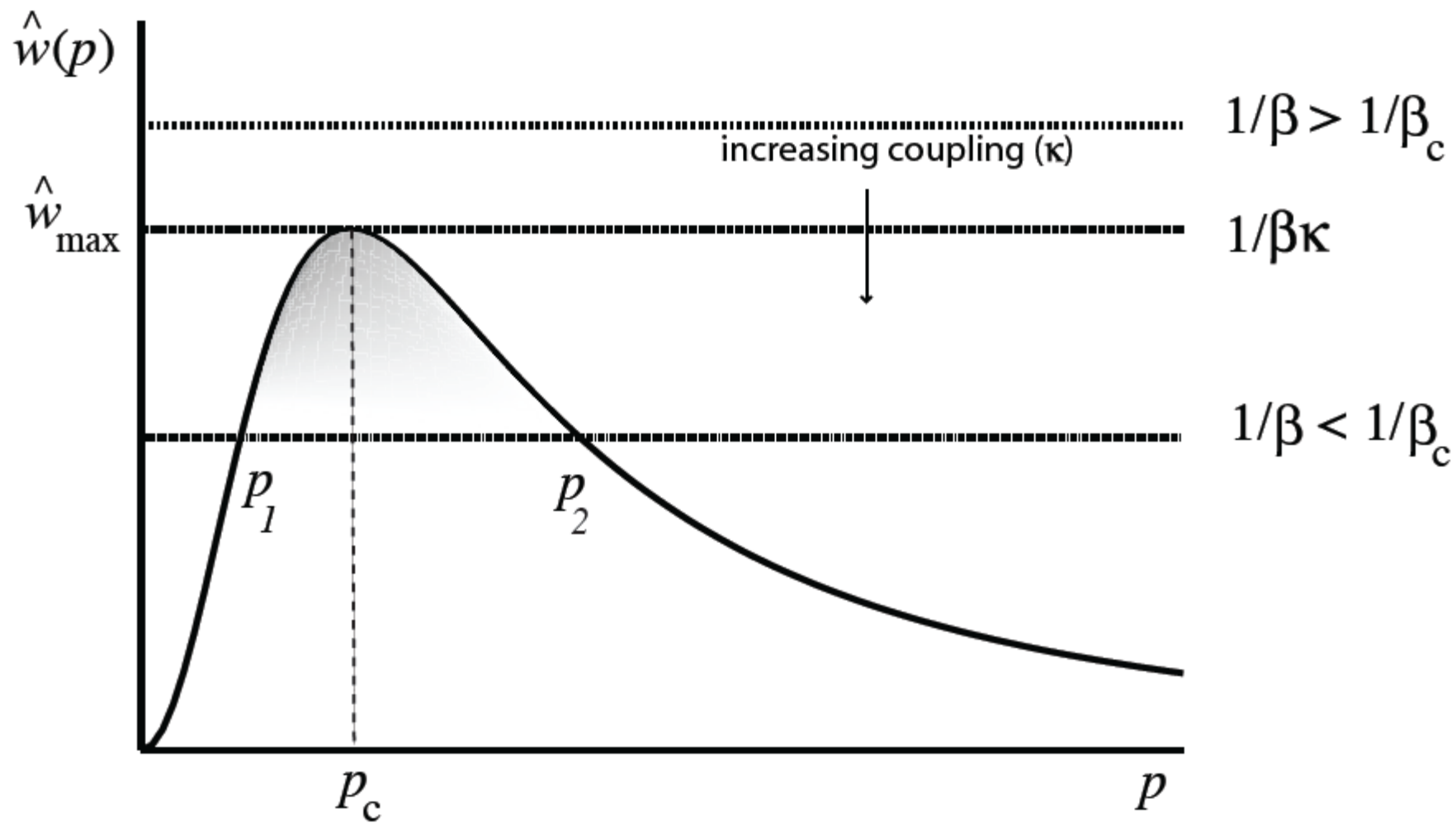
Choosing $\eta(t) = \alpha e^{-\alpha t} H(t)$, we get $\tilde{\eta}(t) = (1 + \lambda/\alpha)^{-1}$,

so that $1 = \kappa\beta(1 + \lambda/\alpha)^{-1}\hat{w}(p)$,

$$\lambda = \alpha(\hat{w}(p)\kappa\beta - 1),$$

$$\hat{w}(p)\kappa\beta < 1,$$

$$\hat{w}(p) < \frac{1}{\kappa\beta}.$$



- For $\beta < \beta_c$ we have $\kappa\hat{w}(p) \leq \kappa\hat{w}_{\max} < 1/\beta$ for all p and the resting state is linearly stable.
- At the critical point $\beta = \beta_c$ we have $\beta_c\kappa\hat{w}(p_c) = 1$ and $\beta_c\kappa\hat{w}(p) < 1$ for all $p \neq p_c$. Hence, $\lambda(p) < 0$ for all $p \neq p_c$, but $\lambda(p_c) = 0$. This signals the point of a *static* instability due to excitation of the pattern $e^{\pm ip_c x}$.
- Beyond the bifurcation point, $\beta > \beta_c$, $\lambda(p_c) > 0$ and this pattern grows with time. In fact there will typically exist a range of values of $p \in (p_1, p_2)$ for which $\lambda(p) > 0$, signalling a set of growing patterns. As the patterns grow, the linear approximation breaks down and nonlinear terms dominate behaviour.
- The saturating property of $f(u)$ tends to create patterns with finite amplitude, that scale as $\sqrt{\beta - \beta_c}$ close to bifurcation and have wavelength $2\pi/p_c$.
- If $p_c = 0$ then we would have a *bulk instability* resulting in the formation of another homogeneous state.

Example: Mexican Hat Function

Biologically motivated choice for $w(x)$:

$$w(x) = \Lambda \left[e^{-\gamma_1|x|} - \Gamma e^{-\gamma_2|x|} \right].$$

$\Lambda = 1$, short-range excitation and long-range inhibition,

$\Lambda = -1$, short-range inhibition and long-range excitation.

$$\hat{w}(p) = 2\Lambda \left[\frac{\gamma_1}{\gamma_1^2 + p^2} - \Gamma \frac{\gamma_2}{\gamma_2^2 + p^2} \right],$$

$$p_c^2 = \frac{\gamma_1^2 \sqrt{\Gamma \gamma_2 / \gamma_1} - \gamma_2^2}{1 - \sqrt{\Gamma \gamma_2 / \gamma_1}}.$$

Full Model

$$h(x, t) = \kappa \int_{-\infty}^t ds F(t - s) \int_{-\infty}^s ds' \eta(s - s') \int_{-\infty}^{\infty} dy w(|x - y|) f(h(y, s' - D(x, y)/v)),$$

with $D(x, y) = |x - y|$.

The homogenous steady state $h(x, t) = h_0$ is

$$h_0 = \kappa f(h_0) \int_0^{\infty} F(s) ds \int_{-\infty}^{\infty} dy w(|y|), \text{ so that we obtain}$$

$$1 = \kappa \beta \hat{w}(p, \lambda) \tilde{\eta}(\lambda) \tilde{F}(\lambda),$$

$$\hat{w}(p, \lambda) = \int_{-\infty}^{\infty} dy w(|y|) e^{-ipy} e^{-\lambda|y|/v}, \quad \beta = f'(h_0).$$

Full Model

In the limit $v \rightarrow \infty$, $\hat{w}(p, \lambda) \rightarrow \hat{w}(p)$.

Then, for $\eta(t) = \alpha e^{-\alpha t} H(t)$, we have

$$1 + \lambda/\alpha = \kappa\beta\hat{w}(p)\tilde{F}(\lambda).$$

For a dynamic instability to occur, we must have

$\text{Re}(\lambda) = 0$ and $\text{Im}(\lambda) \neq 0$, i.e. there must be a pair $\omega, p \neq 0$ s.t. $\lambda = i\omega$ and

$$1 + i\omega/\alpha = \kappa\beta\hat{w}(p)\tilde{F}(i\omega).$$

Full Model

Defining $C(\omega) = \operatorname{Re}(\tilde{F}(i\omega))$, $S(\omega) = \operatorname{Im}(\tilde{F}(i\omega))$, where

$$C(\omega) = \int_0^\infty ds F(s) \cos(\omega s) \leq |C(0)|.$$

Equating the real and imaginary parts, we get

$$1 = \kappa\beta\hat{w}(p)C(\omega), \quad \omega/\alpha = \kappa\beta\hat{w}(p)S(\omega).$$

Dividing the second equation by the first gives

$$\frac{\omega}{\alpha} = \mathcal{H}(\omega), \quad \text{with } \mathcal{H}(\omega) := \frac{S(\omega)}{C(\omega)}.$$

Instability Analysis

Bifurcation condition $\beta = \beta_d$ for a dynamic instability

$$\beta_d \kappa \hat{w}(p_{\min}) = \frac{1}{C(\omega_c)}.$$

Bifurcation condition $\beta = \beta_s$ for a static instability

$$\beta_s \kappa \hat{w}(p_{\max}) = \frac{1}{C(0)}.$$

Assuming $\hat{w}(p_{\min}) < 0 < \hat{w}(p_{\max})$,

dynamic Turing instability if $\beta < \beta_s$ and $p_{\min} \neq 0$,

static Turing instability if $\beta_s < \beta$ and $p_{\max} \neq 0$.

Instability Analysis

Mexican hat with $\Lambda = 1$:

No Turing instability because $p_{\min} = 0$.

Bulk oscillations instead of static patterns when

$$\hat{w}(p_c) < -\frac{C(\omega_c)}{C(0)} |\hat{w}(0)|, \quad p_c^2 = \frac{\gamma_1^2 \sqrt{\Gamma \gamma_2 / \gamma_1} - \gamma_2^2}{1 - \sqrt{\Gamma \gamma_2 / \gamma_1}}.$$

Mexican hat with $\Lambda = -1$:

$p_{\min} = p_c$ and $p_{\max} = 0$. Turing instability when

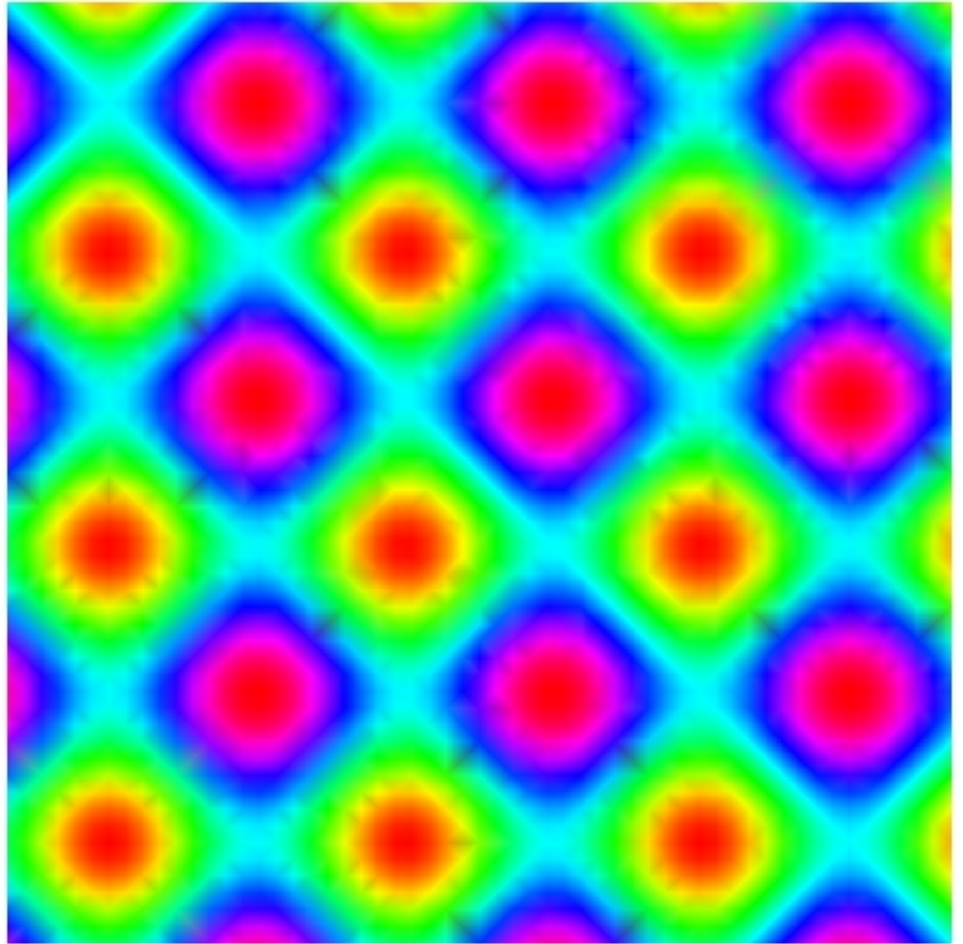
$$\hat{w}(0) < -\frac{C(\omega_c)}{C(0)} |\hat{w}(p_c)|.$$

Doubly Periodic Square Function

$$h(\mathbf{r}) = \sum_j A_j e^{ip_c \mathbf{R}_j \cdot \mathbf{r}}$$

$$\mathbf{R}_1 = \begin{bmatrix} 1 \\ 0 \end{bmatrix}, \mathbf{R}_2 = \begin{bmatrix} 0 \\ 1 \end{bmatrix},$$

$$A_1 \in [0, 2], A_2 = 1.$$



Doubly Periodic Hexagonal Function

$$h(\mathbf{r}) = \sum_j A_j e^{ip_c \mathbf{R}_j \cdot \mathbf{r}}$$

$$\mathbf{R}_1 = \begin{bmatrix} 1 \\ 0 \end{bmatrix},$$

$$\mathbf{R}_2 = \frac{1}{2} \begin{bmatrix} -1 \\ \sqrt{3} \end{bmatrix},$$

$$\mathbf{R}_3 = \frac{1}{2} \begin{bmatrix} 1 \\ \sqrt{3} \end{bmatrix}.$$

