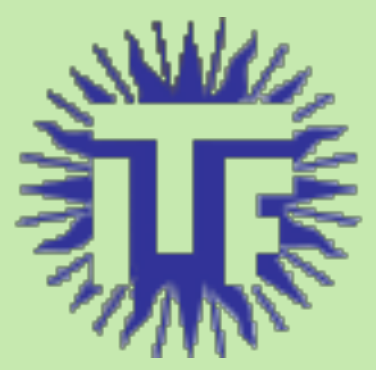


# Towards a semi-holographic superfluid



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## MOTIVATION

**Goal:** using holography describe bosonic *single-particle* spectra in a strongly interacting superfluid!

**What do we mean by single particles?**

- Creation/annihilation operators in Fock space
- Obey canonical commutation relations and the sum rule

$$[\psi_\alpha(\mathbf{x}, t), \psi_{\alpha'}^\dagger(\mathbf{x}', t)]_\pm = \delta(\mathbf{x} - \mathbf{x}')\delta_{\alpha, \alpha'} \quad \stackrel{=}{=} \frac{1}{\pi} \int d\omega \text{Im} G_{R, \alpha, \alpha'}(\omega, \mathbf{k}) = \delta_{\alpha, \alpha'}$$

**Why are they interesting?**

Their spectral functions can be experimentally observed!

## FROM 'STANDARD' HOLOGRAPHY...

- Add a scalar field  $\phi$  to the bulk
- EOM near the boundary yields

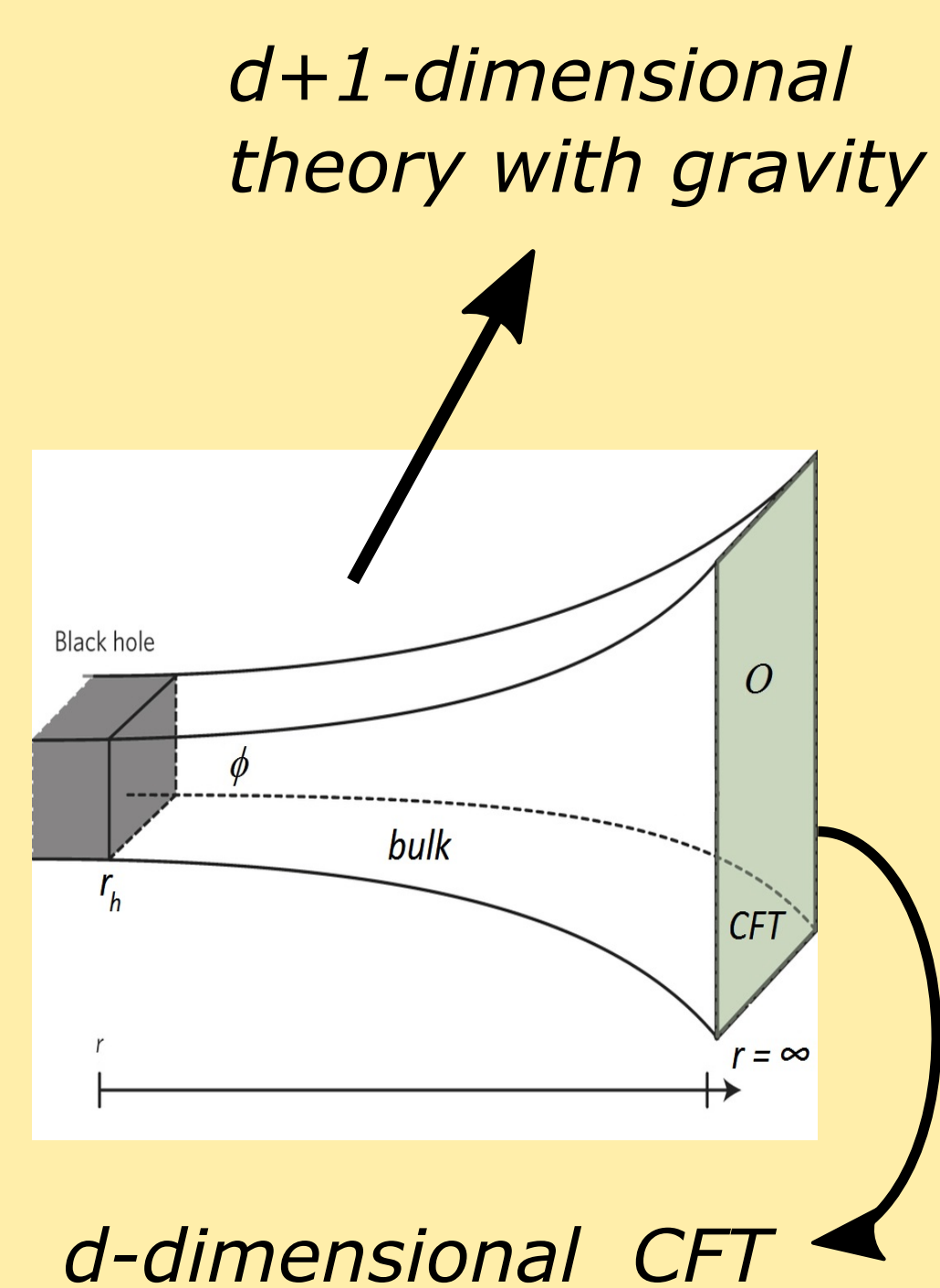
$$\phi \sim \phi_s r^{-\Delta_-} + \phi_O r^{-\Delta_+}$$

source  $\sim$  vev of operator  $O$

- Retarded Green's function of  $O$  behaves as

$$G_{R, O} = (-\omega^2 + k^2 c^2)^{\frac{1}{2}(\Delta_+ - \Delta_-)}$$

which violates the sum rule. Hence  $\phi$  is not an elementary scalar field!



## ...TO SEMI-HOLOGY!

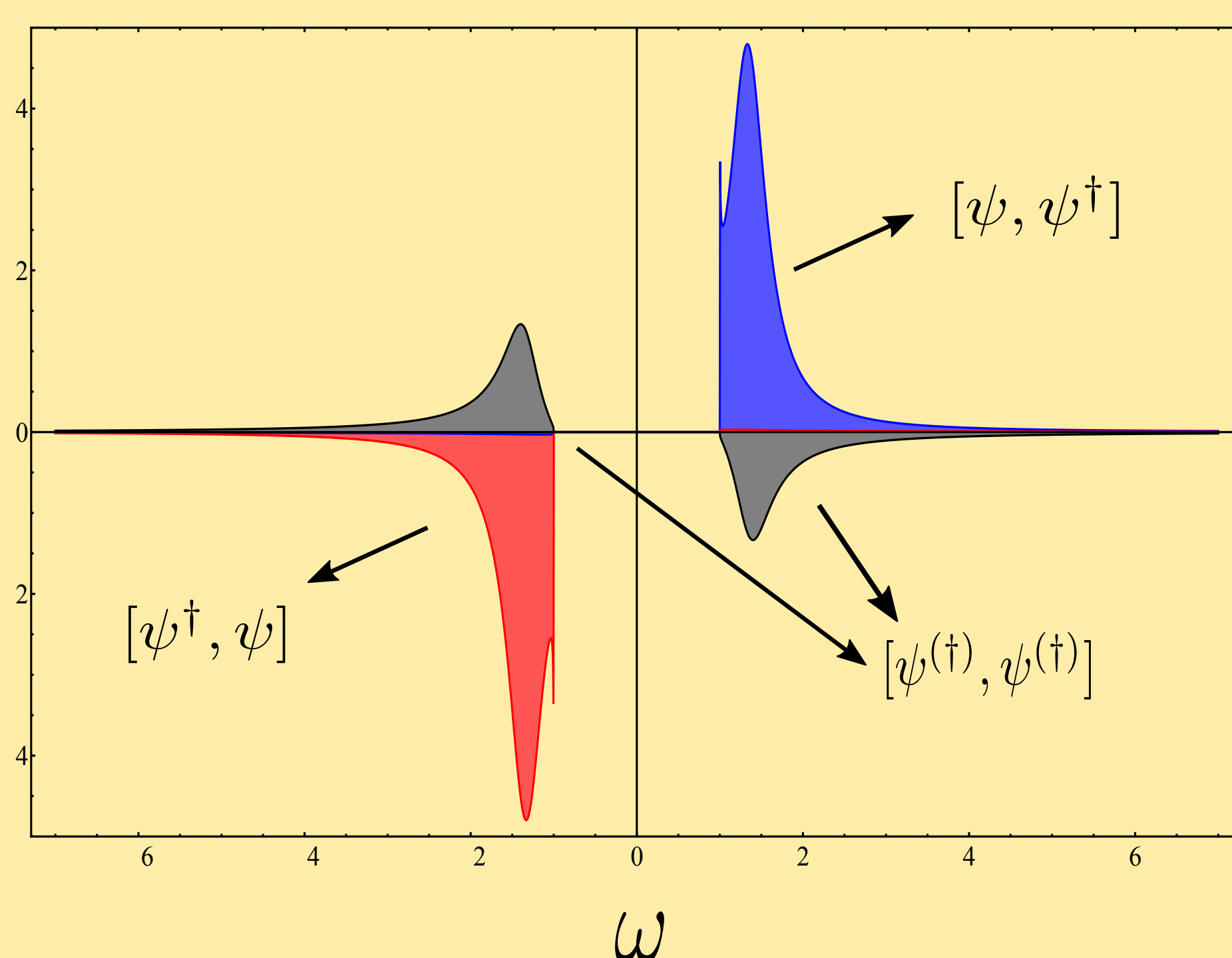
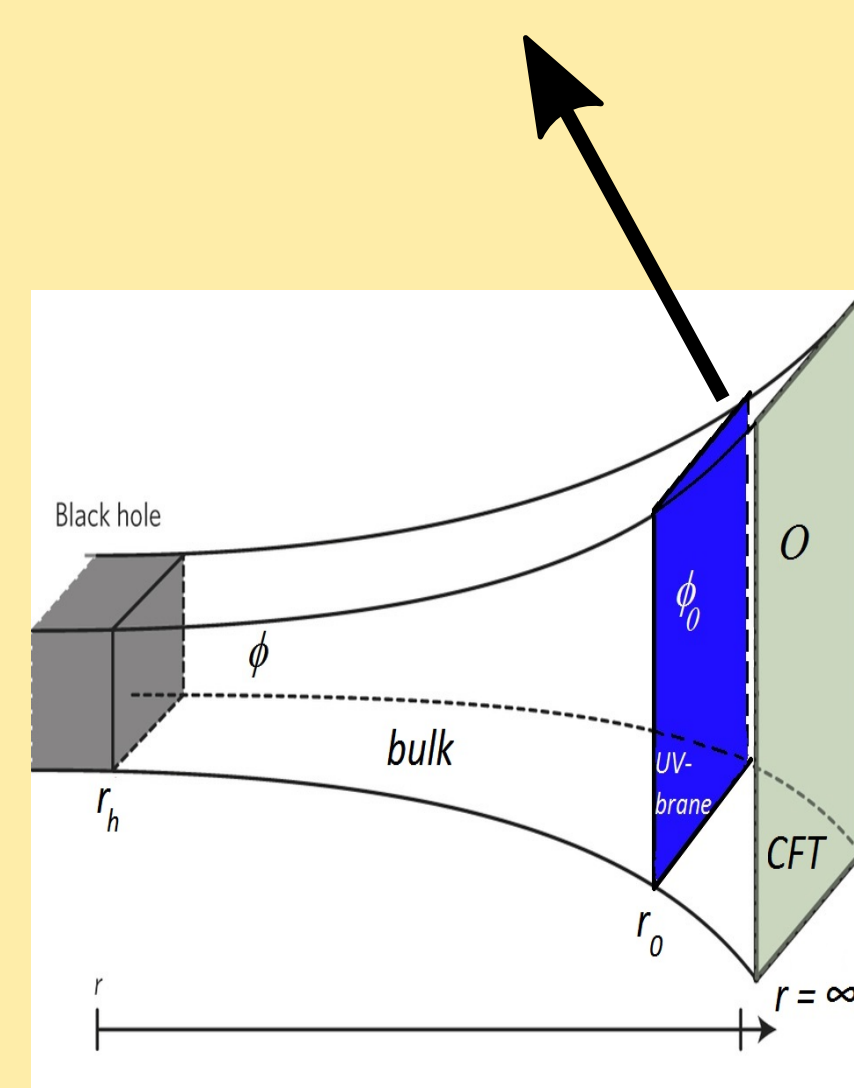
- Introduce a UV brane  $r = r_0$
- Make the source dynamical on this brane [1]!

$$S_{UV} = -\hbar Z \int_{r_0} d^d x \sqrt{-\hbar} \partial_a \phi_s \partial^a \phi_s^*$$

- Retarded Green's function of  $\phi_s$  becomes

$$G_{R, \phi_s} = \frac{-1}{p^2 + g G_{R, O}}$$

UV brane with dynamical source



- $\phi_s$  is an elementary scalar field!

- $G_{R, O}$  now appears as a **self-energy** due to the interaction between  $\phi_s$  and the CFT!

## THE HOLOGRAPHIC SUPERFLUID

The holographic superfluid is a well-known [2] solution to the action

$$S = \int d^5 x \sqrt{-g} \left( \frac{c^3}{16\pi G} (R - \Lambda) - \frac{1}{4\mu_0 c} F_{\mu\nu} F^{\mu\nu} - \hbar g_b \left( D_\mu \phi (D^\mu \phi)^* + \frac{m^2 c^2}{\hbar^2} |\phi|^2 \right) \right)$$

**Properties:**

- Solution consists of metric, temporal gauge field and scalar hair
- Above some temperature  $T_c$ ,  $\phi = 0$  only solution => left with a charged black hole
- Below  $T_c$  there are solutions with  $\phi_s = 0$  and  $\phi_O \neq 0$ . On the boundary this yields spontaneous symmetry breaking!

## EFFECTIVE ACTION

To find the semi-holographic Green's function, we must first understand the holographic result.

Near  $T_c$ , one can describe a *complex* order parameter  $\Delta$  with the following effective Landau-Ginzburg action

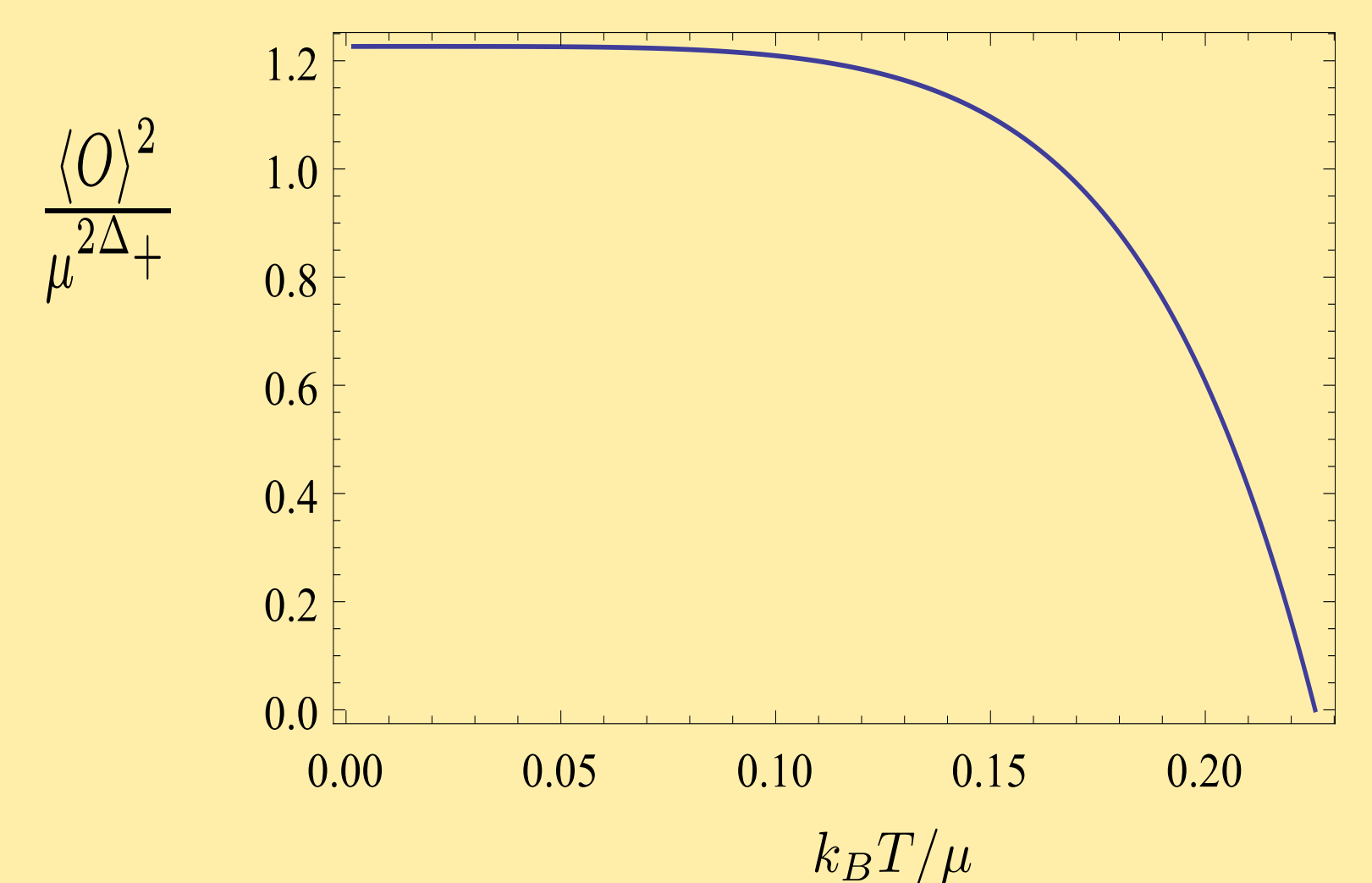
$$S = \int d^4 x \left( -i a \Delta^* \partial_t \Delta + \gamma |\nabla \Delta|^2 + \alpha_0 (T - T_c) |\Delta|^2 + \frac{\beta}{2} |\Delta|^4 \right)$$

**Relationship between  $\Delta$  and  $O$**

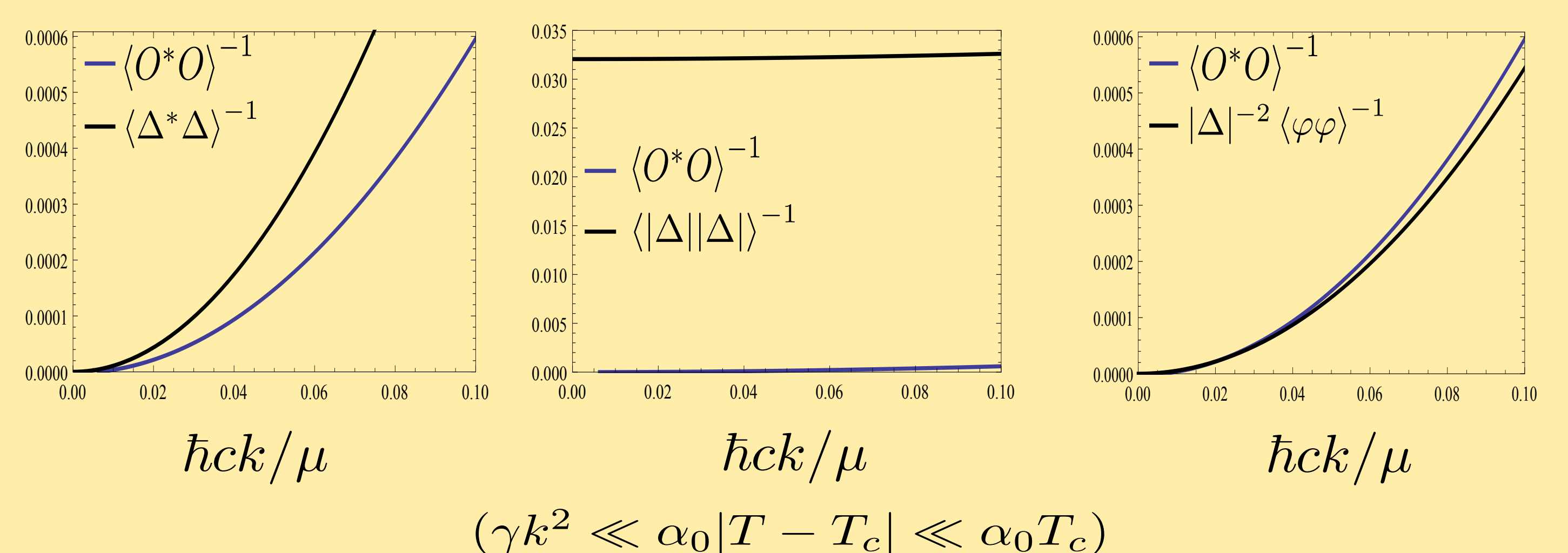
- Below  $T_c$  we find

$$\langle \Delta \rangle = \sqrt{\frac{-\alpha_0 (T - T_c)}{\beta}}$$

which agrees with  $\langle O \rangle(T)$ .



- Assuming  $O \sim \Delta$  one can determine the above coefficients from the retarded Green's function of  $O$ . See [3] for  $T > T_c$ !
- The Green's functions below  $T_c$  do not behave as expected from  $O \sim \Delta$ . Rather, they seem to describe **phase** fluctuations! We believe this is due to large- $N$  arguments.



## REFERENCES

- [1] U. Gürsoy, E. Plauschinn, H. Stoof and S. Vandoren, JHEP 1205, 018 (2012)
- [2] S. Hartnoll, C. Herzog and G. Horowitz, JHEP 0812, 015 (2008)
- [3] J.-H. She, B. Overbosch, Y.-W. Sun, Y. Liu, K. Schalm, J. Mydosh and J. Zaanen, Phys.Rev. B84, 144527 (2011)



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